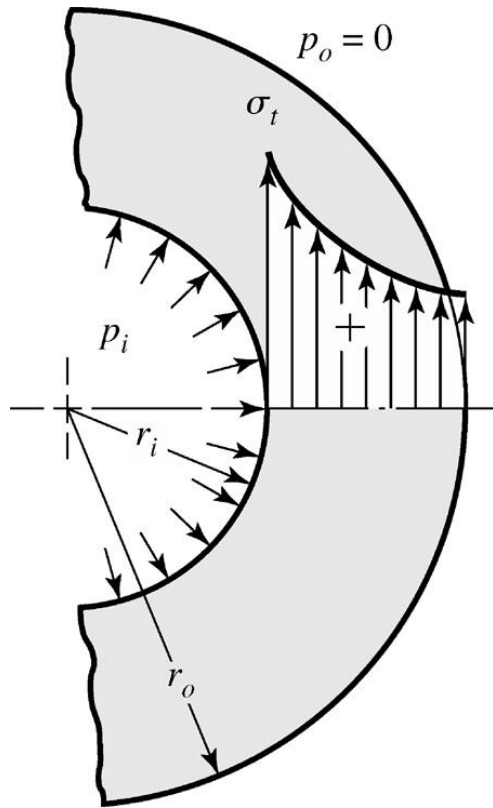
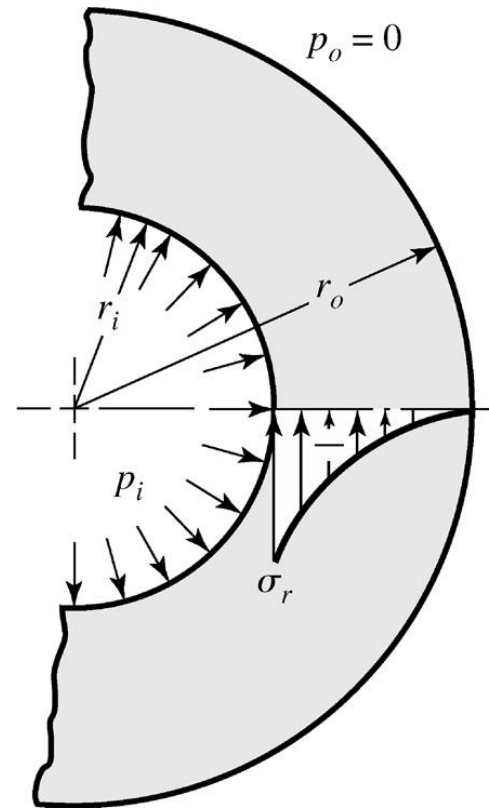


# Stress Analysis II

# Stress distributions in a thick walled cylinder subjected to internal pressure



(a) Tangential stress distribution




(b) Radial stress distribution

# Thermal Stresses and Strains

$$\varepsilon = \varepsilon_T + \varepsilon_E \quad \varepsilon_T : \text{Thermal strain} \quad \varepsilon_E : \text{Elastic Strain}$$

Uniform temperature change:

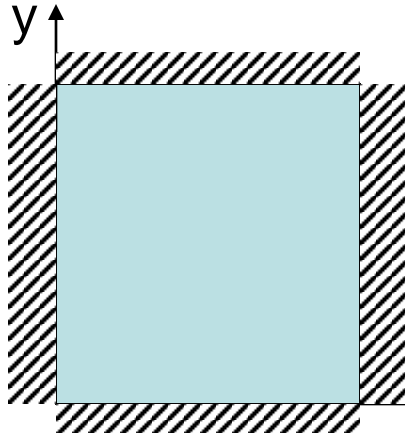
if there is no constraint, there is no stress.



$$\varepsilon_x = (\varepsilon_x)_T + (\varepsilon_x)_E = 0 \quad (\varepsilon_x)_E = -(\varepsilon_x)_T$$

$$\sigma_x = E (\varepsilon_x)_E = -E (\varepsilon_x)_T = -E \alpha \Delta T$$

Plane Stress:



$$(\varepsilon_x)_T = (\varepsilon_y)_T = \alpha \Delta T \quad \varepsilon_x = (\varepsilon_x)_T + (\varepsilon_x)_E = 0$$

$$(\varepsilon_x)_E = -(\varepsilon_x)_T = -\alpha \Delta T \quad \text{but} \quad (\varepsilon_y)_E = (\varepsilon_x)_E$$

$$\sigma_x = \frac{E}{1-\nu^2} [(\varepsilon_x)_E + \nu(\varepsilon_y)_E]$$

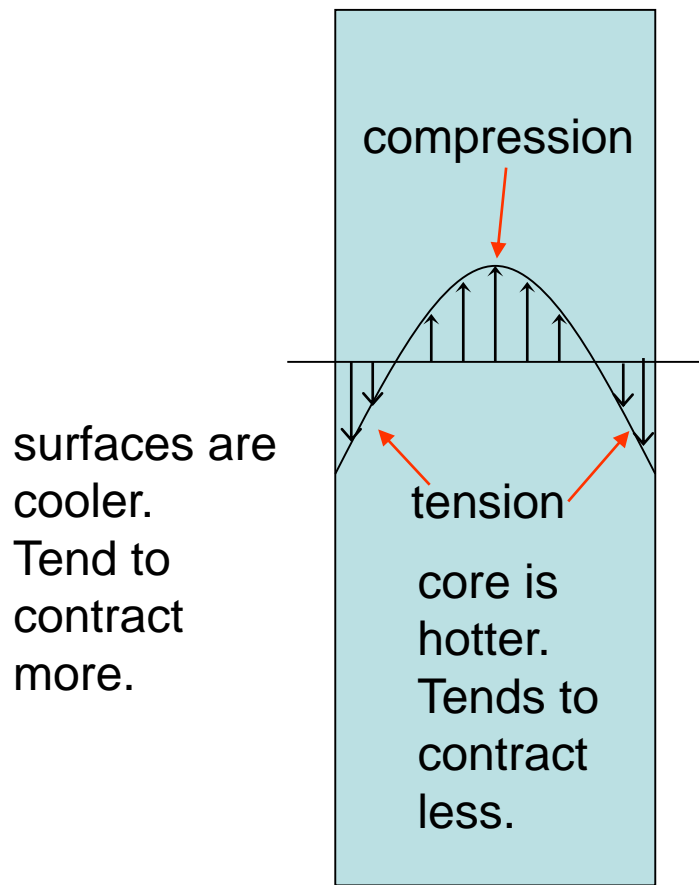
$$\sigma_x = \frac{E(1+\nu)(\varepsilon_x)_E}{1-\nu^2}$$

$$\sigma_x = -\frac{E\alpha\Delta T}{1-\nu}$$

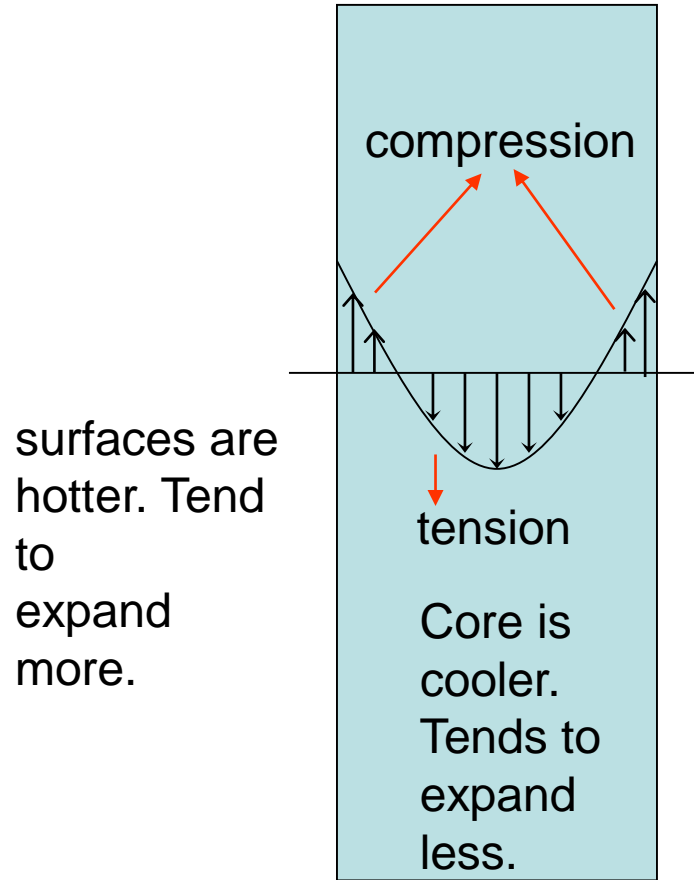
temperature stress arise because of constraint.

THERMAL STRESSES arise because of temperature gradient in a member.

Infinite slab during heating and cooling:



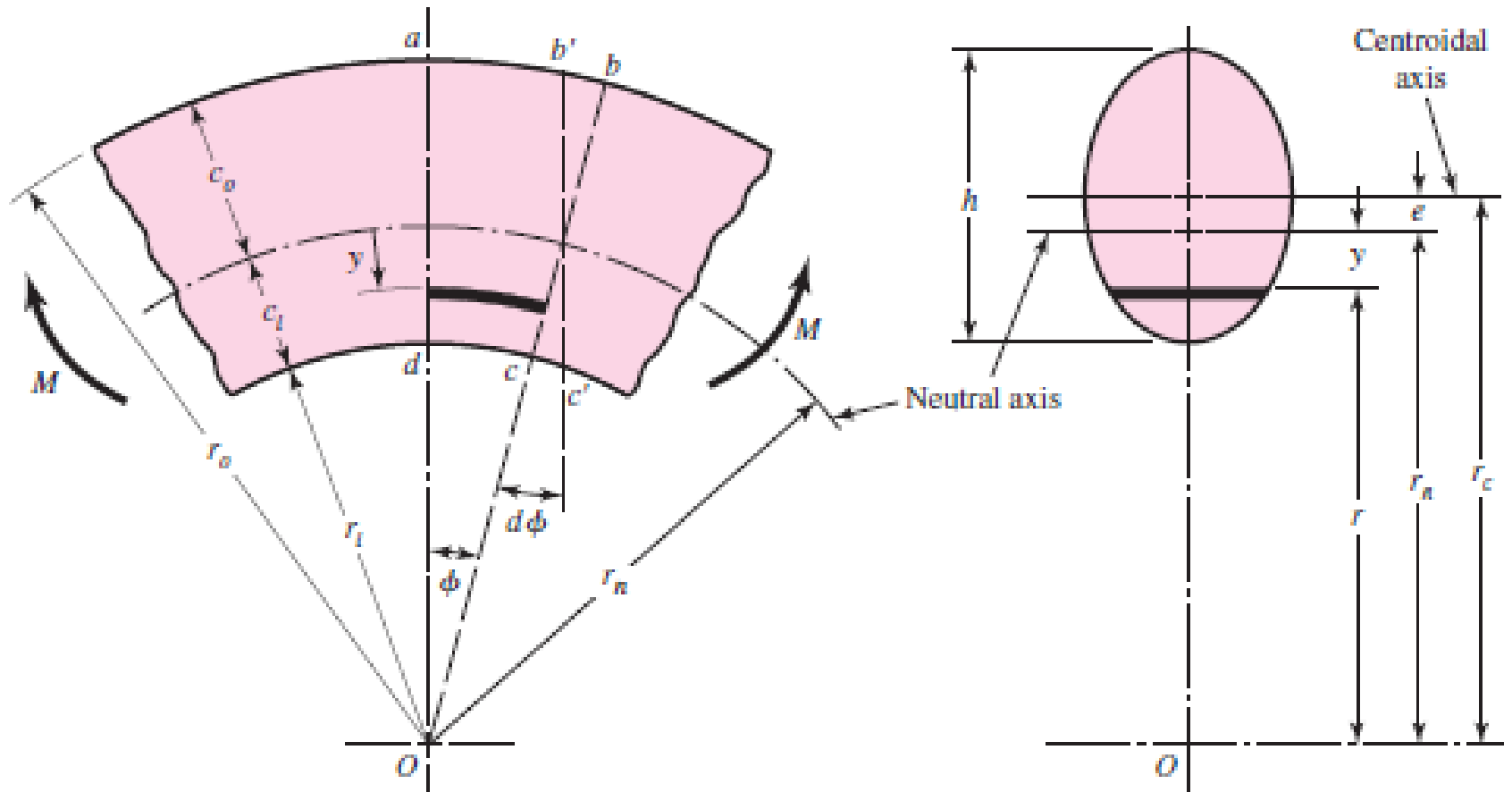
Cooling



Heating

# Curved Members in Flexure

- Distribution of stress in a curved flexural member is determined by using the following assumptions.
  - Cross-section has an axis of symmetry.
  - Plane cross-sections remain plane after bending.
  - Modulus of elasticity is the same in tension and compression



For curved beams;  
 neutral axis is not coincident with centroidal axis, and  
 stress distribution is not linear.

bc rotates through  $d\Phi$  to b'c'.

$$\text{Strain on a fiber at } r; \varepsilon = \frac{\delta}{\ell} = \frac{(r_n - r)d\phi}{r\phi} = \frac{yd\phi}{r\phi} \dots\dots\dots(a)$$

$$\text{Corresponding normal stress; } \sigma = E\varepsilon = \frac{E\delta}{\ell} = \frac{E(r_n - r)d\phi}{r\phi} = \frac{Eyd\phi}{r\phi} \dots\dots\dots(b)$$

There is no external force on the beam, therefore the sum of the normal forces on the cross-section of the beam must be zero.

$$\int_A \sigma dA = E \frac{d\phi}{\phi} \int_A \frac{(r_n - r)}{r} dA = 0 \dots\dots\dots(c)$$

$$E \frac{d\phi}{\phi} \left( r_n \int_A \frac{dA}{r} - \int_A dA \right) = 0 \dots\dots\dots(d)$$

$$r_n \int_A \frac{dA}{r} - A = 0 \dots\dots\dots(e)$$

$$r_n = \frac{A}{\int_A \frac{dA}{r}}$$

location of neutral axis

Note that location of centroid is given by  
hence  $r_c$  and  $r_n$  are not the same.

$$r_c = \frac{\int_A r dA}{A}$$

The external moment on the beam is equal to the moment of the normal forces on the cross-section of the beam.

$$M = \int_A y \sigma dA = \int_A (r_n - r) \sigma dA = E \frac{d\phi}{\phi} \int_A \frac{(r_n - r)^2}{r} dA \quad \dots\dots\dots(f)$$

note that  $(r_n - r)^2 = r_n^2 - 2r_n r + r^2$

$$M = E \frac{d\phi}{\phi} \left( \underbrace{r_n^2 \int_A \frac{dA}{r} - r_n \int_A dA - r_n \int_A dA + \int_A r dA}_{\dots\dots\dots(g)} \right)$$

$$r_n \left( r_n \int_A \frac{dA}{r} - \int_A dA \right) = 0$$

$$M = E \frac{d\phi}{\phi} \left( - \underbrace{r_n \int_A dA}_A + \underbrace{\int_A r dA}_{r_c A} \right)$$

$$M = E \frac{d\phi}{\phi} (r_c - r_n) A = E \frac{d\phi}{\phi} e A$$



from (b)

$$\frac{d\phi}{\phi} = \frac{r\sigma}{E(r_n - r)}$$

$$M = E \frac{r\sigma}{E(r_n - r)} eA$$

$$\sigma = \frac{M(r_n - r)}{reA}$$

but  $y = r_n - r$

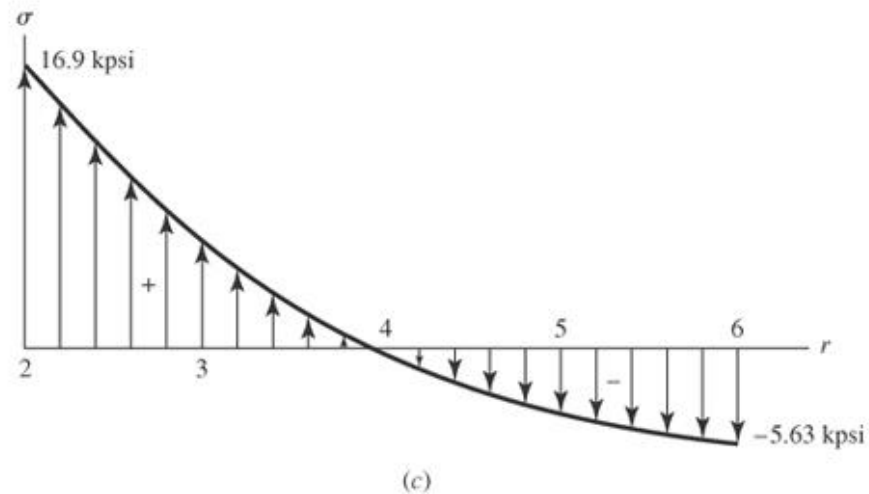
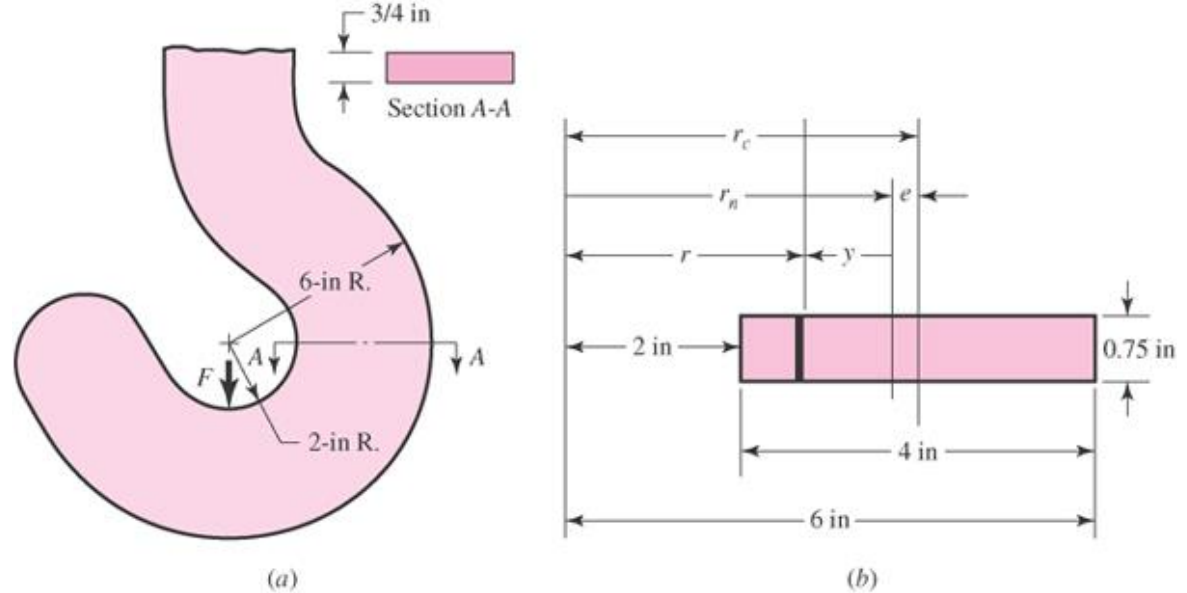
$$r = r_n - y$$

$$\sigma = \frac{My}{(r_n - y)eA}$$

$$\sigma_i = \frac{Mc_i}{r_i eA} \quad \sigma_o = -\frac{Mc_o}{r_o eA}$$

If the bending moment  $M$  is due to a force  $F$ , the moment is taken about centroidal axis rather than neutral axis. Also, the direct normal stress  $\pm F/A$  should be superposed with bending stress.

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In curved beam problems, one needs to take the integrals;

$$r_n = \frac{A}{\int_A \frac{dA}{r}} \text{ and}$$

$$r_c = \frac{\int_A r dA}{A}$$

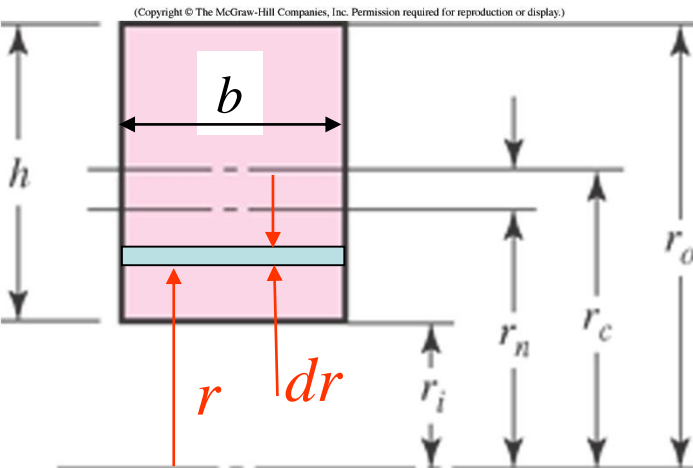
For some other common cross sections (Circular, trapezoidal, I, box) the integrals are listed in a table, in the text book.)

Optimum cross section:

The cross section is most efficiently utilized when the maximum tensile and compressive stresses are equal.

Then for the case of pure bending,

$$\sigma_i = \sigma_o \quad \text{and} \quad \frac{c_i}{c_o} = \frac{r_i}{r_o}$$



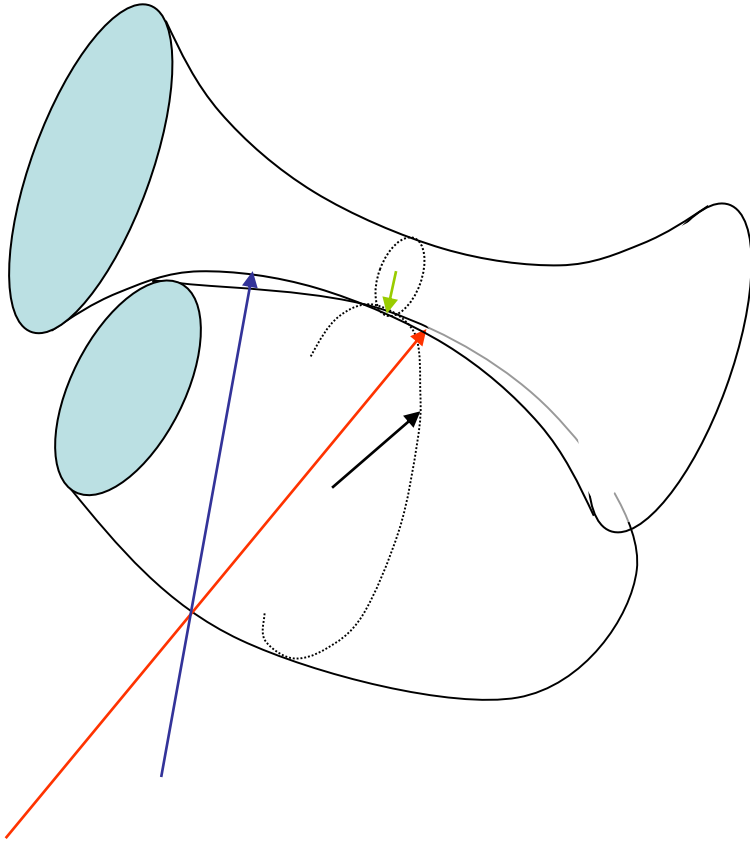
$$r_n = \frac{A}{\int_A \frac{dA}{r}} = \frac{bh}{\int_A \frac{bdr}{r}} = \frac{h}{\ln \frac{r_o}{r_i}}$$

$$r_c = r_i + h/2$$

# Contact Stresses

- When two bodies having curved surfaces are pressed together;
  - point or line contact changes to area contact
  - 3D stress state develops in the vicinity of contact area (local stresses)
  - stresses may cause failure such as cracking, pitting, flaking.
- Examples: wheel on rail, cam-follower, pin in a bearing, mating gear teeth, rollers on raceway of an anti-friction bearing.

Most general case occurs when both bodies have double radius of curvature.



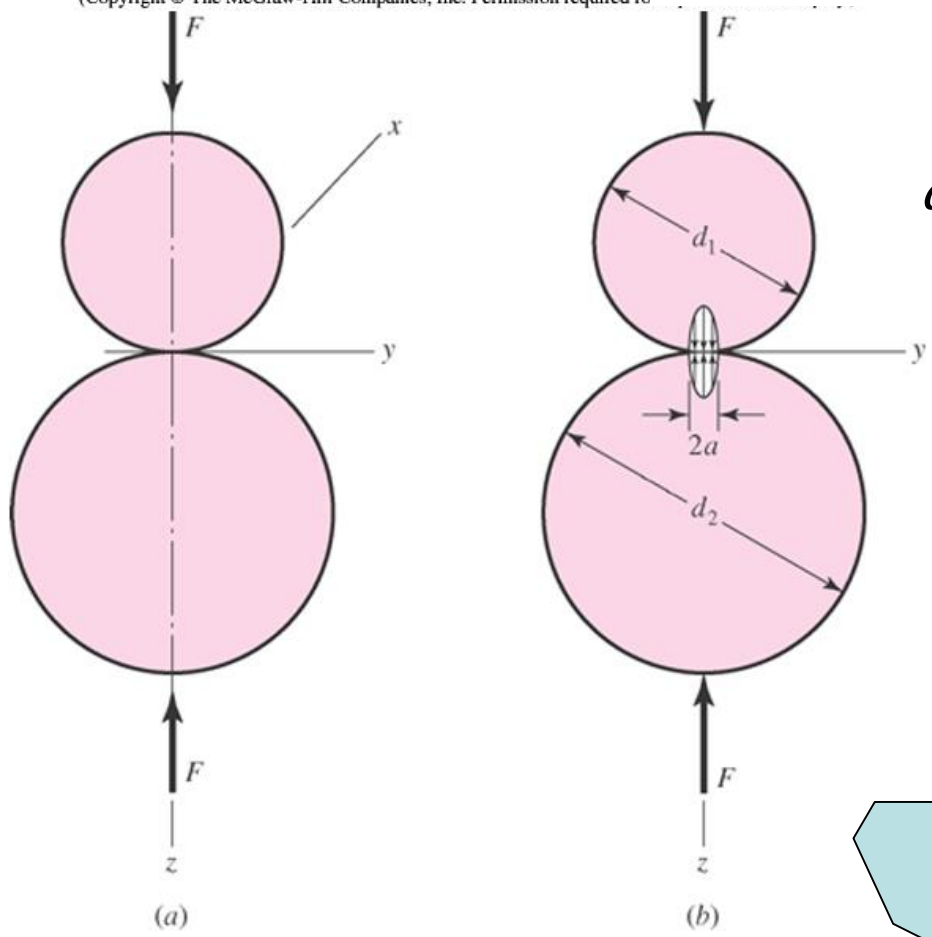
We consider only two special cases of practical importance:

- i) Contacting spheres
- ii) contacting cylinders

The results presented here are due to Hertz, so the contact stresses are also called Hertzian stresses.

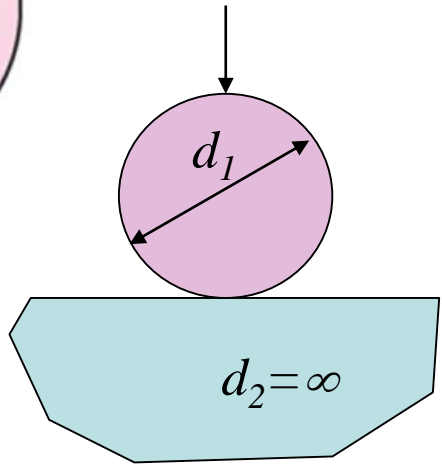
Contacting Spheres: Hemispherical pressure distribution is formed. Max. pressure is at the center of the circular contact area, of diameter  $2a$ .

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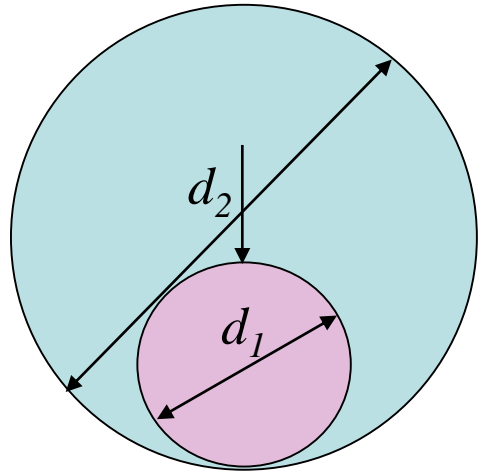


$$a = \sqrt[3]{\frac{3F}{8} \frac{(1-\nu_1^2)/E_1 + (1-\nu_2^2)/E_2}{1/d_1 + 1/d_2}}$$

$$P_{\max} = \frac{3F}{2\pi a^2}$$



$$1/d_2 = 0$$



$$d_2 < 0$$

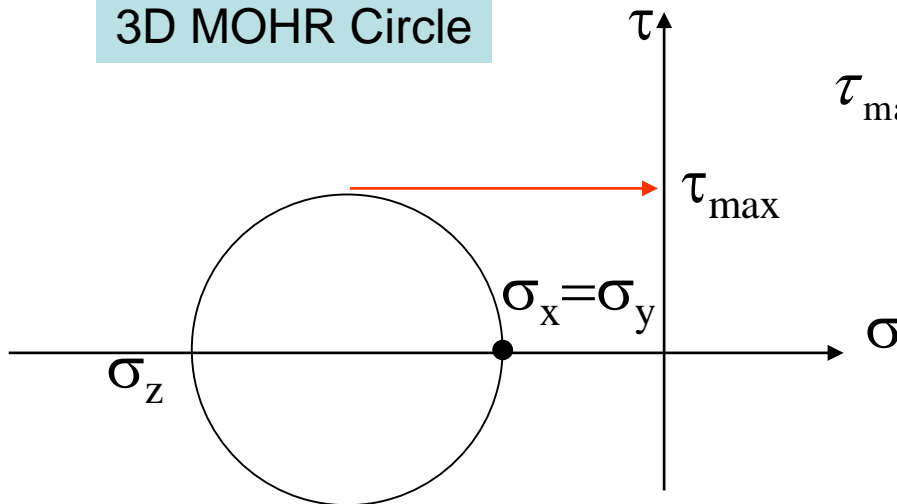
Maximum stresses occur on z-axis and these are principal stresses.

$$\sigma_x = \sigma_y = -p_{\max} \left[ (1 + \nu) \left( 1 - \frac{|z|}{a} \tan^{-1} \left( \frac{1}{|z/a|} \right) \right) - \frac{1}{2(1 + z^2/a^2)} \right]$$

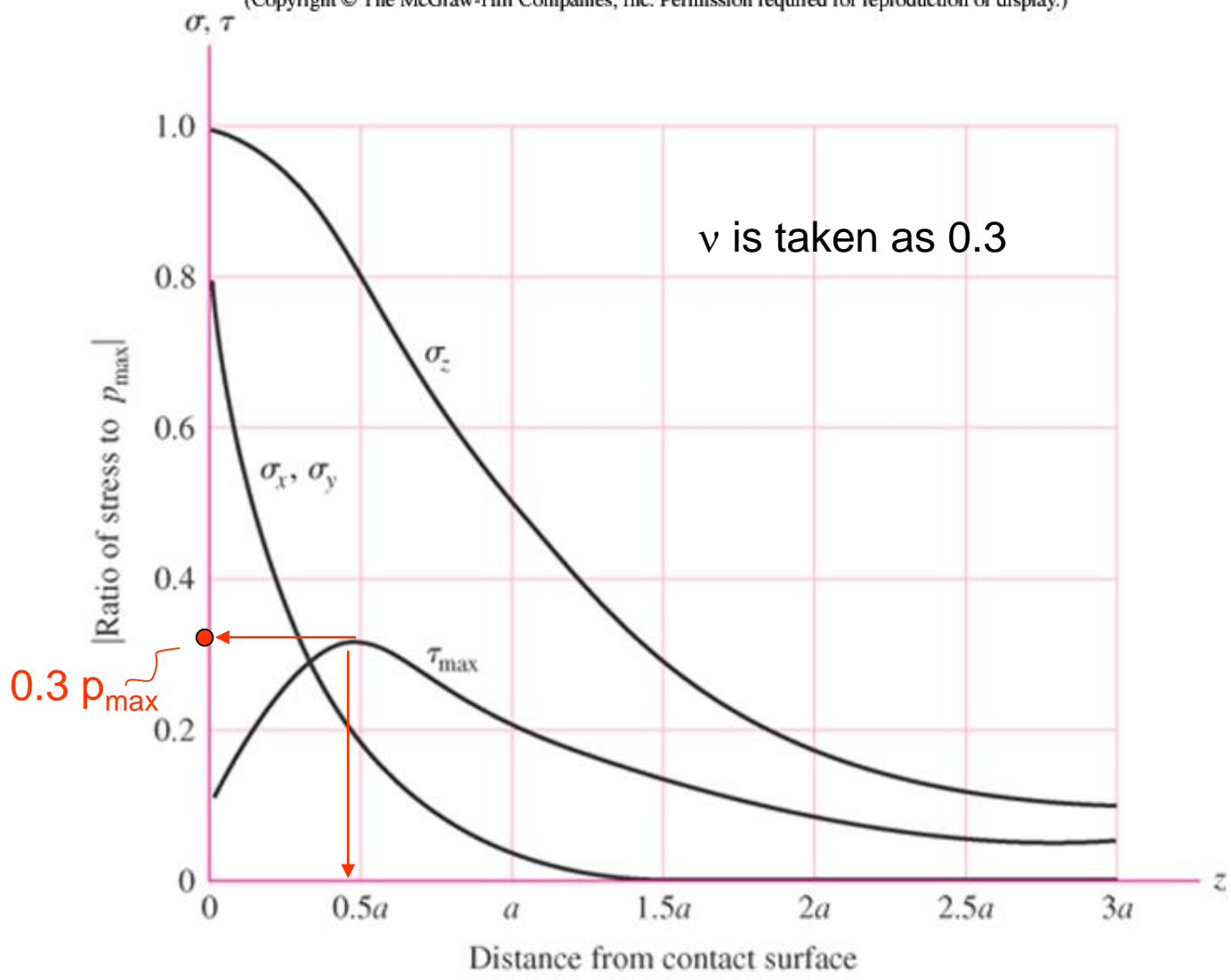
$$\sigma_z = -\frac{p_{\max}}{1 + z^2/a^2}$$

we considered stresses only on z-axis (x=y=0) for design and analysis purposes, since the maximum stresses and the critical point (the point which is most likely to fail) are on z-axis.

### 3D MOHR Circle



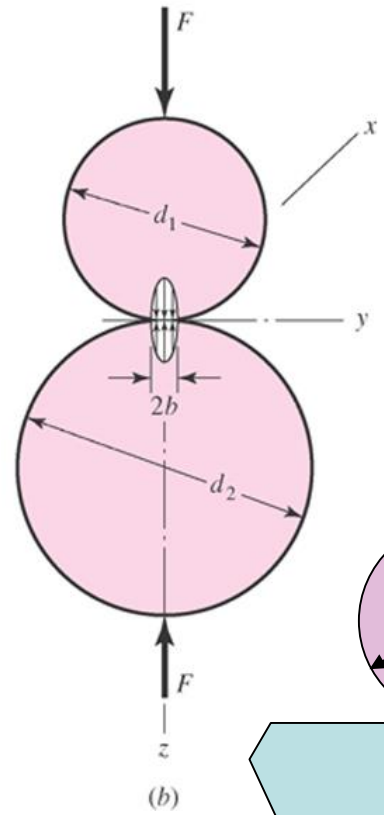
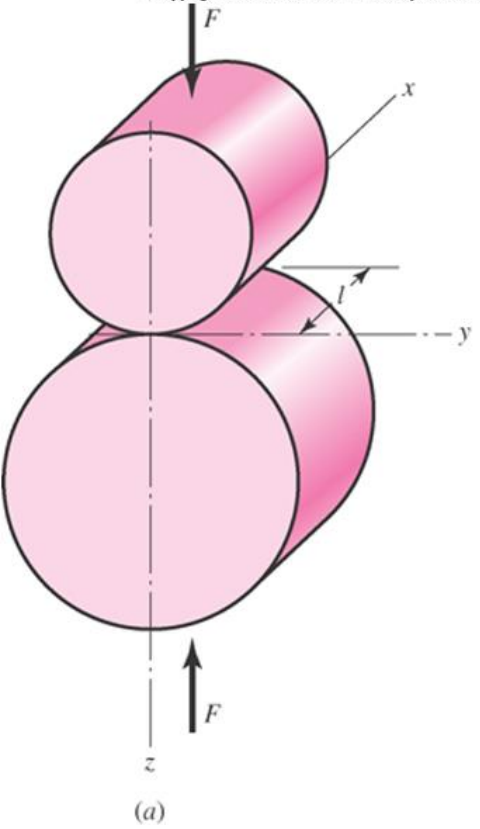
$$\tau_{\max} = \frac{\sigma_x - \sigma_z}{2} = \frac{\sigma_y - \sigma_z}{2}$$



The crack originates at the location of maximum shear stress, which is slightly below the surface, and propagates towards the surface.

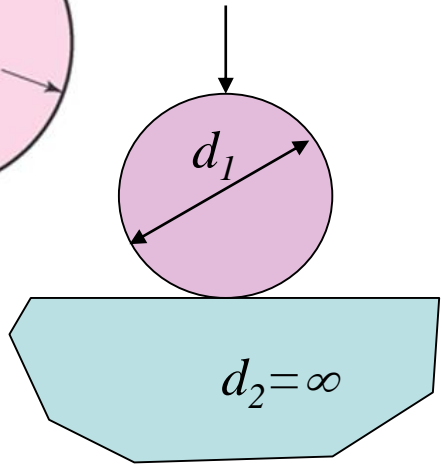
Contacting Cylinders: Elliptical pressure distribution is formed. Max. pressure is at  $y=0$ .

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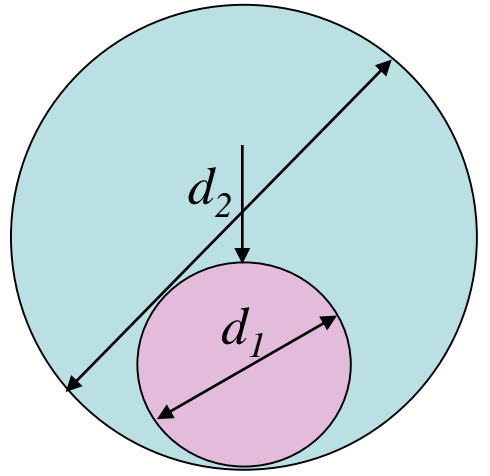


$$b = \sqrt{\frac{2F}{\pi l} \frac{(1-\nu_1^2)/E_1 + (1-\nu_2^2)/E_2}{1/d_1 + 1/d_2}}$$

$$P_{\max} = \frac{2F}{\pi b l}$$



$$1/d_2 = 0$$



$$d_2 < 0$$



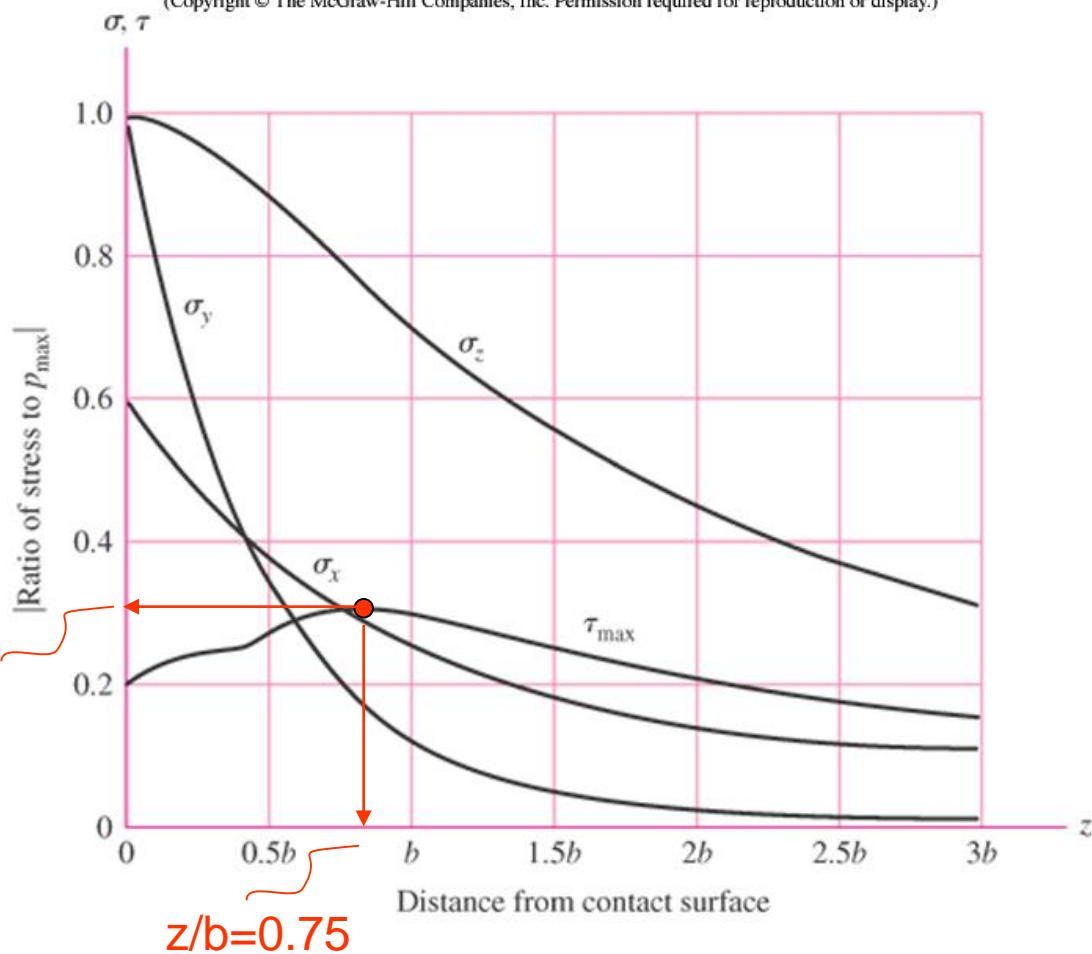
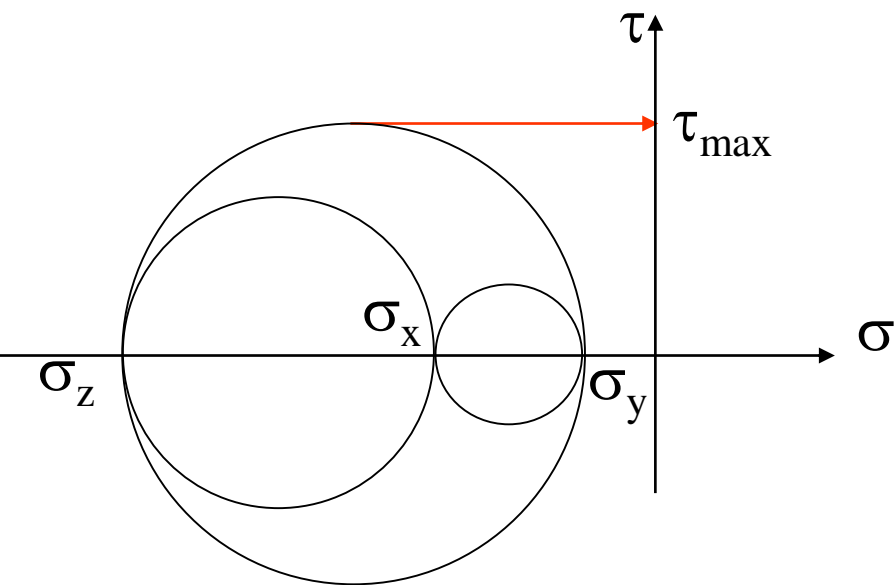
# Stresses along z-axis :

$$\sigma_x = -2\nu p_{\max} \left[ \sqrt{1 + \frac{z^2}{b^2}} - \left| \frac{z}{b} \right| \right]$$

$$\sigma_y = -p_{\max} \left[ \frac{1 + 2(z^2/b^2)}{\sqrt{1 + z^2/b^2}} - 2 \left| \frac{z}{b} \right| \right]$$

$$\sigma_z = -\frac{p_{\max}}{\sqrt{1 + z^2/b^2}}$$

## 3D MOHR Circle



$$\tau_{\max} = \frac{\sigma_z - \sigma_y}{2}$$